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# Computational simulation of complex fine-scale heterogeneous flows using homogenization and coarse-graining approaches

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Joint work with Nathan March (QUT) and Ian Turner (QUT)

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- ▶ Fine-scale diffusion model:

$$\frac{\partial u}{\partial t} + \nabla \cdot (-D(\mathbf{x})\nabla u) = 0, \quad x \in \Omega \subset \mathbb{R}^2.$$

- ▶ If the scale at which the diffusivity  $D(\mathbf{x})$  changes is small compared to the size of the domain  $\Omega$ , then the amount of computation required to solve this model is prohibitive due to the very fine mesh required to capture the heterogeneity.
- ▶ This can be overcome by homogenizing or partially-homogenizing the heterogeneous medium  $\Omega$ .
- ▶ Homogenized diffusion model:

$$\frac{\partial U}{\partial t} + \nabla \cdot (-\mathbf{D}_{\text{eff}}(\mathbf{x})\nabla U) = 0, \quad x \in \Omega \subset \mathbb{R}^2,$$

where  $U(\mathbf{x}, t)$  is a smoothed/coarse-scale approximation to the fine-scale solution  $u(\mathbf{x}, t)$ .

- ▶ Cell problem for first column of  $\mathbf{D}_{\text{eff}}$ :

$$\nabla \cdot (D(\mathbf{x})\nabla(\psi + x)) = 0, \quad \mathbf{x} = (x, y) \in Y = [0, L]^2,$$

$$\psi(\mathbf{x}) \text{ is periodic with period } Y, \quad \frac{1}{L^2} \int_Y \psi \, dV = 0,$$

$$\mathbf{D}_{\text{eff}}(:, 1) = \frac{1}{L^2} \int_Y D(\mathbf{x})\nabla(\psi + x) \, dV.$$

- ▶ Cell problem for second column of  $\mathbf{D}_{\text{eff}}$ :

$$\nabla \cdot (D(\mathbf{x})\nabla(\psi + y)) = 0, \quad \mathbf{x} = (x, y) \in Y = [0, L]^2,$$

$$\psi(\mathbf{x}) \text{ is periodic with period } Y, \quad \frac{1}{L^2} \int_Y \psi \, dV = 0,$$

$$\mathbf{D}_{\text{eff}}(:, 2) = \frac{1}{L^2} \int_Y D(\mathbf{x})\nabla(\psi + y) \, dV.$$

- ▶ Definitions involving non-periodic boundary conditions also exist (e.g. confined or uniform boundary conditions).

- ▶ For a layered medium, the cell problems can be solved exactly yielding the following effective diffusivity:

$$D(\mathbf{x}) = \begin{cases} D_A & \text{if } 0 < y < \frac{L}{2}, \\ D_B & \text{if } \frac{L}{2} < y < L, \end{cases} \implies \mathbf{D}_{\text{eff}} = \begin{bmatrix} D_a & 0 \\ 0 & D_h \end{bmatrix},$$

where  $D_a$  and  $D_h$  are the arithmetic and harmonic means:

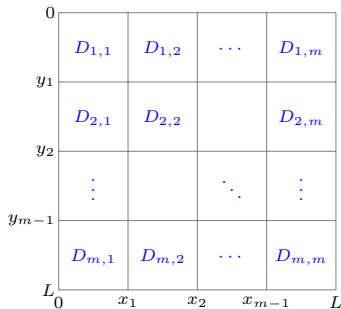
$$D_a = \frac{D_A + D_B}{2}, \quad D_h = \frac{2D_A D_B}{D_A + D_B}.$$

- ▶ For complex geometries, numerical methods are required (Carr and Turner, 2014; Rupp et al., 2018; Szymkiewicz and Lewandowska, 2006).
- ▶ The goal of this work is to develop a semi-analytical method for solving the cell problems and computing  $\mathbf{D}_{\text{eff}}$ .

# Semi-analytical solution

## Block heterogeneous medium

- ▶ Complex heterogeneous geometries can be represented as an array of blocks.
- ▶ Consider the  $Y = [0, L]^2$  consisting of an  $m^2$  grid of rectangular blocks:



- ▶ Each block is isotropic with its own diffusivity value.
- ▶ Consider cell problem for  $\mathbf{D}_{\text{eff}}(:, 1)$  (second column follows similarly).

- ▶ On each block domain, the cell problem is expressed as

$$0 = \nabla \cdot (D_{i,j} \nabla (\psi_{i,j} + x)),$$

where  $D_{i,j}$  is the diffusivity in the  $(i, j)$ th block (row  $i$ , column  $j$ ).

- ▶ For a block heterogeneous domain (discontinuous diffusivity) we also require boundary conditions at the interfaces between adjacent blocks.
- ▶ Solution and the flux are continuous at each interface:

- Horizontal interfaces:

$$\psi_{i,j} = \psi_{i+1,j}, \quad D_{i,j} \frac{\partial \psi_{i,j}}{\partial y} = D_{i+1,j} \frac{\partial \psi_{i+1,j}}{\partial y}.$$

- Vertical interfaces:

$$\psi_{i,j} = \psi_{i,j+1}, \quad D_{i,j} \left( \frac{\partial \psi_{i,j}}{\partial x} + 1 \right) = D_{i,j+1} \left( \frac{\partial \psi_{i,j+1}}{\partial x} + 1 \right).$$

- ▶ Change of variable:

$$v_{i,j} = \psi_{i,j} + x.$$

- ▶ On each block domain, the cell problem is now expressed as

$$\nabla^2 v_{i,j} = 0.$$

- ▶ Solution and the flux are continuous at each interface:

- Horizontal interfaces:

$$v_{i,j} = v_{i+1,j}, \quad D_{i,j} \frac{\partial v_{i,j}}{\partial y} = D_{i+1,j} \frac{\partial v_{i+1,j}}{\partial y}.$$

- Vertical interfaces:

$$v_{i,j} = v_{i,j+1}, \quad D_{i,j} \frac{\partial v_{i,j}}{\partial x} = D_{i+1,j} \frac{\partial v_{i+1,j}}{\partial x}.$$

- ▶ Introduce unknown functions for the diffusive fluxes at interfaces between adjacent blocks, e.g.,

$$D_{i,j} \frac{\partial v_{i,j}}{\partial x} = D_{i,j+1} \frac{\partial v_{i,j+1}}{\partial x} =: g_{(j-1)m+i+1}(y), \quad (\text{vertical interfaces}).$$

$$D_{i,j} \frac{\partial v_{i,j}}{\partial y} = D_{i+1,j} \frac{\partial v_{i+1,j}}{\partial y} =: q_{(j-1)m+i+1}(x), \quad (\text{horizontal interfaces}).$$

- ▶ This allows us to consider the following general block problem:

$$\nabla^2 v_{i,j} = 0,$$

$$D_{i,j} \frac{\partial v_{i,j}}{\partial x} = g_{(i-1)n+j}(y), \quad \text{at } x = x_{j-1}, \quad D_{i,j} \frac{\partial v_{i,j}}{\partial x} = g_{(i-1)n+j+1}(y), \quad \text{at } x = x_j,$$

$$D_{i,j} \frac{\partial v_{i,j}}{\partial y} = q_{(j-1)m+i}(x), \quad \text{at } y = y_{i-1}, \quad D_{i,j} \frac{\partial v_{i,j}}{\partial y} = q_{(j-1)m+i+1}(x), \quad \text{at } y = y_i.$$



- Solution on each block:

$$\begin{aligned}v_{i,j}(x, y) = & -\frac{a_{i,j,0}}{4l_j}(x - x_j)^2 + \frac{b_{i,j,0}}{4l_j}(x - x_{j-1})^2 - \frac{c_{i,j,0}}{4h_i}(y - y_i)^2 + \frac{d_{i,j,0}}{4h_i}(y - y_{i-1})^2 \\ & - h_i \sum_{k=1}^{\infty} \frac{a_{i,j,k}}{\gamma_{i,j,k}} \cosh \left[ \frac{k\pi(x - x_j)}{h_i} \right] \cos \left[ \frac{k\pi(y - y_{i-1})}{h_i} \right] \\ & + h_i \sum_{k=1}^{\infty} \frac{b_{i,j,k}}{\gamma_{i,j,k}} \cosh \left[ \frac{k\pi(x - x_{j-1})}{h_i} \right] \cos \left[ \frac{k\pi(y - y_{i-1})}{h_i} \right] \\ & - l_j \sum_{k=1}^{\infty} \frac{c_{i,j,k}}{\mu_{i,j,k}} \cosh \left[ \frac{k\pi(y - y_i)}{l_j} \right] \cos \left[ \frac{k\pi(x - x_{j-1})}{l_j} \right] \\ & + l_j \sum_{k=1}^{\infty} \frac{d_{i,j,k}}{\mu_{i,j,k}} \cosh \left[ \frac{k\pi(y - y_{i-1})}{l_j} \right] \cos \left[ \frac{k\pi(x - x_{j-1})}{l_j} \right] + K_{i,j},\end{aligned}$$

where  $\gamma_{i,j,k} = k\pi \sinh \frac{k\pi l_j}{h_i}$  and  $\mu_{i,j,k} = k\pi \sinh \frac{k\pi h_i}{l_j}$ ,  $h_i = y_i - y_{i-1}$  and  $l_j = x_j - x_{j-1}$ .

- ▶ Coefficients are integrals of unknown flux functions, e.g.

$$a_{i,j,k} = \frac{2}{h_i} \int_{y_{i-1}}^{y_i} \frac{g^{(i-1)n+j}(y)}{D_{i,j}} \cos\left(\frac{k\pi(y - y_{i-1})}{h_i}\right) dy.$$

- ▶ We approximate these integrals numerically using a midpoint rule, e.g.

$$a_{i,j,k} \approx \frac{2}{D_{i,j}h_i} \sum_{p=1}^N \omega_p g^{(i-1)n+j}(y_p) \cos\left(\frac{k\pi(y_p - y_{i-1})}{h_i}\right),$$

where  $N$  is the number of abscissas per interface and  $\omega_p$  and  $y_p$  are the appropriate weights and abscissas.

- ▶ Quadrature approximation requires the evaluations of the unknown interface functions at the abscissas, e.g.  $g^{(i-1)n+j}(y_p)$ .
- ▶ By determining these evaluations, we can compute the coefficients (e.g.  $a_{i,j,k}$ ) and thus compute the effective diffusivity.

- ▶ Enforce continuity of the solution at the abscissas along each interface, e.g.

$$v_{i+1,j}(x_p, y_i) - v_{i,j}(x_p, y_i) = 0 \quad (\text{horizontal interface}).$$

- ▶ This yields a system of linear equations that can be solved for the evaluations of the unknown interface functions:

$$\mathbf{Ax} = \mathbf{b},$$

where  $\mathbf{x}$  is a vector of dimension  $m^2(N+1)$  containing the required evaluations.

- ▶ As we have an analytical expression for the solution of the interface functions, the entries of  $\mathbf{D}_{\text{eff}}$  can be expressed in terms of the coefficients, e.g.

$$\mathbf{D}_{\text{eff}}(1, 1) = \frac{1}{L^2} \sum_{i=1}^m \sum_{j=1}^n \left[ \frac{D_{i,j} A_{i,j} (a_{i,j,0} + b_{i,j,0})}{4} + l_j^2 \sum_{k=1}^{\infty} \frac{(c_{i,j,k} - d_{i,j,k}) [1 - (-1)^k]}{k\pi} \right],$$

where  $A_{i,j} = l_j h_i$  is the area of the  $(i, j)$ th block.

# Linear system dimension

## Comparison to a standard numerical method

- ▶  $m$  by  $m$  array of square blocks.
- ▶  $N$  abscissas per interface.
- ▶ Assume spacing between abscissas and nodes is equal.
- ▶ Linear system:

$$\mathbf{Ax} = \mathbf{b}$$

- ▶ Finite volume method:  
Dimension of  $\mathbf{x}$ :  $m^2N^2$ .
- ▶ Semi-analytical method:  
Dimension of  $\mathbf{x}$ :  $m^2(2N + 1)$ .

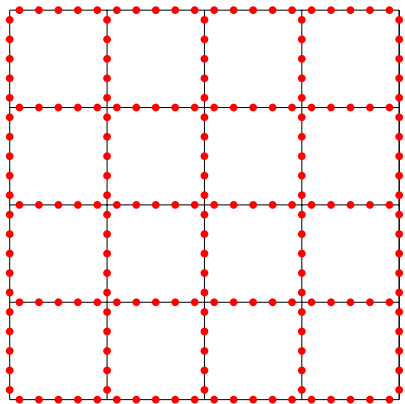


Figure 1: Abscissas ( $4 \times 4$  array of blocks).

# Linear system dimension

## Comparison to a standard numerical method

- ▶  $m$  by  $m$  array of square blocks.
- ▶  $N$  abscissas per interface.
- ▶ Assume spacing between abscissas and nodes is equal.
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Dimension of  $\mathbf{x}$ :  $m^2 N^2$ .
- ▶ Semi-analytical method:  
Dimension of  $\mathbf{x}$ :  $m^2(2N + 1)$ .

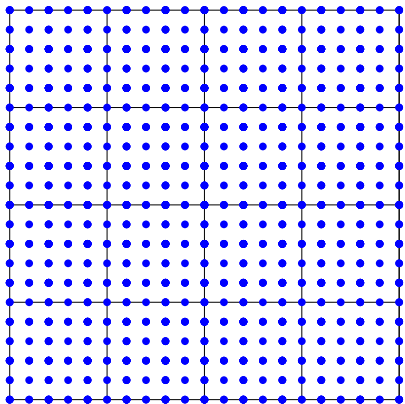


Figure 2: Nodes ( $4 \times 4$  array of blocks).



Figure 3: Standard test case (Szymkiewicz, 2013):  $4 \times 4$  array of blocks (Diffusivity: dark = 0.1, light = 1).

# Semi-analytical method

## Comparison to standard numerical method



N	Semi-Analytical		Finite Volume	
	$ (D_{\text{eff}} - \tilde{D}_{\text{eff}}) / D_{\text{eff}} $	Runtime (s)	$ (D_{\text{eff}} - \tilde{D}_{\text{eff}}) / D_{\text{eff}} $	Runtime (s)
4	$\begin{pmatrix} 6.84\text{e-}3 & 5.04\text{e-}3 \\ 5.04\text{e-}3 & 4.47\text{e-}3 \end{pmatrix}$	0.00747	$\begin{pmatrix} 1.30\text{e-}2 & 2.44\text{e-}3 \\ 2.44\text{e-}3 & 8.47\text{e-}3 \end{pmatrix}$	0.00923
8	$\begin{pmatrix} 3.01\text{e-}3 & 2.21\text{e-}3 \\ 2.21\text{e-}3 & 1.98\text{e-}3 \end{pmatrix}$	0.0109	$\begin{pmatrix} 4.82\text{e-}3 & 1.88\text{e-}3 \\ 1.88\text{e-}3 & 3.14\text{e-}3 \end{pmatrix}$	0.0277
16	$\begin{pmatrix} 1.40\text{e-}3 & 1.02\text{e-}3 \\ 1.02\text{e-}3 & 9.23\text{e-}4 \end{pmatrix}$	0.0331	$\begin{pmatrix} 1.75\text{e-}3 & 9.12\text{e-}4 \\ 9.12\text{e-}4 & 1.14\text{e-}3 \end{pmatrix}$	0.115
32	$\begin{pmatrix} 6.77\text{e-}4 & 4.94\text{e-}4 \\ 4.94\text{e-}4 & 4.48\text{e-}4 \end{pmatrix}$	0.0629	$\begin{pmatrix} 6.17\text{e-}4 & 3.76\text{e-}4 \\ 3.76\text{e-}4 & 4.02\text{e-}4 \end{pmatrix}$	0.530
64	$\begin{pmatrix} 3.42\text{e-}4 & 2.50\text{e-}4 \\ 2.50\text{e-}4 & 2.27\text{e-}4 \end{pmatrix}$	0.270	$\begin{pmatrix} 2.05\text{e-}4 & 1.36\text{e-}4 \\ 1.36\text{e-}4 & 1.33\text{e-}4 \end{pmatrix}$	2.92

$\tilde{D}_{\text{eff}}$ : Approximate  $D_{\text{eff}}$  (semi-analytical or finite volume method)

$D_{\text{eff}}$ : Benchmark  $D_{\text{eff}}$  using finite volume method with a very fine grid.

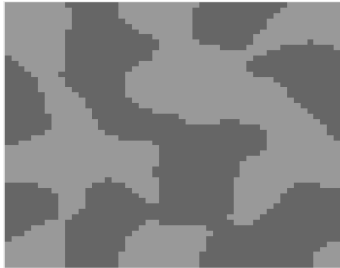
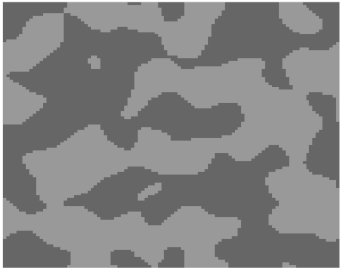
	50 × 50	100 × 100
		
$\mathbf{D}_{\text{eff}}$	$\begin{pmatrix} 0.310 & 0.0177 \\ 0.0177 & 0.342 \end{pmatrix}$	$\begin{pmatrix} 0.340 & 0.000954 \\ 0.000954 & 0.304 \end{pmatrix}$

Table 1: Diffusivity: dark = 0.1, light = 1.



Preprint available on the arXiv repository: <https://arxiv.org/abs/1812.06680>.

## A fast semi-analytical homogenization method for block heterogeneous media

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<sup>b</sup>ARC Centre of Excellence for Mathematical and Statistical Frontiers (ACEMS), Queensland University of Technology (QUT), Brisbane, Australia.

- ▶ New semi-analytical method for solving boundary value problems on block locally-isotropic heterogeneous media.
- ▶ Method provides explicit formula for effective diffusivity  $\mathbf{D}_{\text{eff}}$ .
- ▶ While achieving equivalent accuracy, semi-analytical method is faster than a standard finite volume method for the test problems we considered.
- ▶ Improved efficiency is mainly due to the much smaller linear system size required.

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- Rupp, A., Knabner, P., and Dawson, C. (2018). A local discontinuous Galerkin scheme for Darcy flow with internal jumps. *Computat Geosci*, 22(4):1149–1159.
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